

2004

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Recommended Citation

Glimm, Tilman and Oliker, Vladimir, "Optical Design of Two-reflector Systems, the Monge-Kantorovich Mass Transfer Problem and Fermat's Principle" (2004). *Mathematics*. 58.

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Optical Design of Two-reflector Systems, the Monge-Kantorovich Mass Transfer Problem and Fermat's Principle

TILMANN GLIMM & VLADIMIR OLIKER

ABSTRACT. It is shown that the problem of designing a two-reflector system transforming a plane wave front with given intensity into an output plane front with prescribed output intensity can be formulated and solved as the Monge-Kantorovich mass transfer problem.

1. INTRODUCTION

Consider a two-reflector system of configuration shown schematically on Fig. 1. Let $(x = (x_1, x_2, \dots, x_n), z)$ be the Cartesian coordinates in \mathbb{R}^{n+1} , $n \geq 2$, with z being the horizontal axis and x_1, x_2, \dots, x_n the coordinates in the hyperplane $\alpha : z = 0$. Let B_1 denote a beam of parallel light rays propagating in the positive z -direction and let $\bar{\Omega}$ denote the wavefront which is the cross section of B_1 by hyperplane α . Assume that Ω is a bounded domain on α . An individual ray of the front is labeled by a point $x \in \bar{\Omega}$. The light intensity of the beam B_1 is denoted by $I(x)$, $x \in \Omega$, where I is a non-negative integrable function.

The incoming beam B_1 is intercepted by the first reflector R_1 , defined as a graph of a function $z(x)$, $x \in \bar{\Omega}$. The rays in B_1 are reflected off R_1 forming a beam of rays B_2 . The beam B_2 is intercepted by reflector R_2 , which transforms it

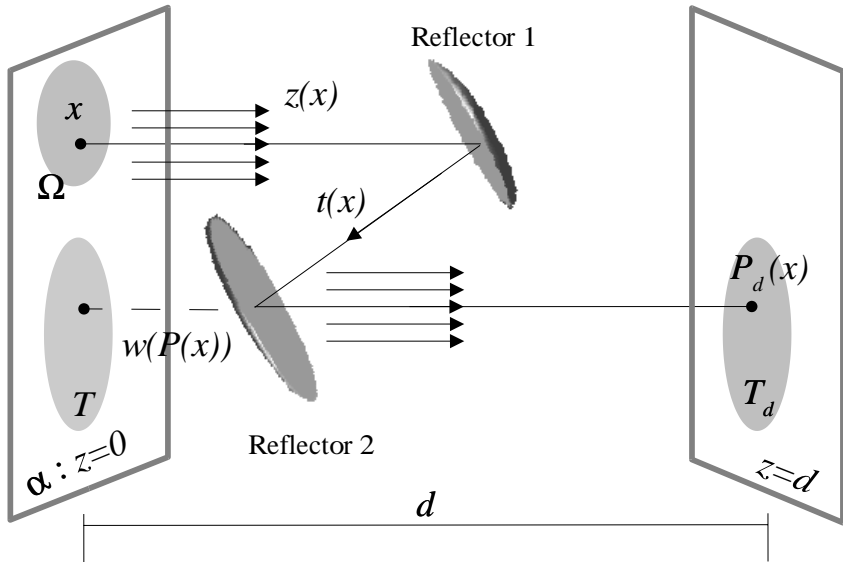


FIGURE 1. Sketch for Problem I

into the output beam B_3 . The beam B_3 also consists of parallel light rays propagating in the same direction as B_1 . The output wavefront at a distance $d > 0$ from the hyperplane α is denoted by \bar{T}_d ; we denote by \bar{T} the projection of \bar{T}_d on the hyperplane α . The second reflector R_2 is also assumed to be a graph of a function $w(p)$, $p \in \bar{T}$. The quantity $1/|J(P_d(x))|$, where P_d is the map of $\bar{\Omega}$ on \bar{T}_d and J is the Jacobian, is the *expansion ratio* and it measures the expansion of a tube of rays due to the two reflections [9]. It is assumed that both R_1 and R_2 are perfect reflectors and no energy is lost in the transformation process. Consequently, the corresponding relation between the input intensity I on Ω and output intensity L on T_d is given by

$$(1.1) \quad L(P_d(x))|J(P_d(x))| = I(x).$$

The “two-reflector” problem that needs to be solved by designers of optical systems consists in determining the reflectors R_1 and R_2 so that all of the properties of the two-reflector system above hold for prescribed in advance domains Ω , T , and positive integrable functions $I(x)$, $x \in \Omega$, and $L(p)$, $p \in T$; see Malyak [10] and other references there. It is usually assumed in applications that Ω and T_d are bounded and convex.

Two fundamental principles of geometrical optics are used to describe the transformation of the beam B_1 into beam B_3 : the classical reflection law leading to the ray tracing equations defining the map P_d , and the energy conservation law

for the energy flux along infinitesimally small tubes of rays; see [10], where the problem is formulated for rotationally symmetric data and a class of rotationally symmetric solutions is found.

The problem of recovering reflectors R_1 and R_2 without assuming rotational symmetry was formulated rigorously by Oliker and Prussner in [11], and it was shown that it can be considered as a problem of determining a special map of $\bar{\Omega} \rightarrow \bar{T}$ with a potential satisfying an equation of Monge-Ampère type relating the input and output intensities. Existence and uniqueness of weak solutions were established by Oliker at that time but only the numerical results implementing a constructive scheme for proving existence were presented in [11] for several test cases. Detailed proofs were given in [13].

In this paper we show that this problem can also be studied in the framework of the Monge-Kantorovich mass transfer problem studied by Brenier [3], Caffarelli [6], Gangbo and McCann [7], and other authors. In our notation, the Monge-Kantorovich mass transfer problem is to transfer the intensity I on Ω into the intensity L on T via a map $P : \Omega \rightarrow T$ for which the total transportation cost $\int_{\Omega} C(x, P(x))I dx$ is minimal. Here $C(x, p)$ is a given strictly convex cost function.

The proof of existence and uniqueness of solutions to the Monge-Kantorovich problem is obtained by solving a minimization problem for the functional

$$(1.2) \quad (\zeta, \omega) \mapsto \int_{\Omega} \zeta I dx - \int_T \omega L dp,$$

considered on pairs of continuous functions ζ on $\bar{\Omega}$ and ω on \bar{T} that satisfy

$$(1.3) \quad \zeta(x) - \omega(p) \geq -C(x, p), \quad x \in \bar{\Omega}, p \in \bar{T}.$$

Under various conditions it is shown in [3], [6], [7] that this functional is minimized by some pair (ζ_0, ω_0) (referred to as Kantorovich potentials), and that $P(x) = x + \nabla \zeta_0$ solves the Monge-Kantorovich problem.

Applying these ideas, we show that the geometric optics problem at hand can be formulated as a Monge-Kantorovich mass transfer problem with a quadratic cost function; see Section 6. The Kantorovich potentials correspond to the pair of reflectors that solve the problem. The condition (1.3) has a geometric meaning; namely, it filters out reflectors that allow only optical paths longer than a certain prescribed one. The functional (1.2) to be minimized is the mean horizontal distance between points of the two-reflectors, with the average weighted by the two intensities.

We prove that there are always two different reflector systems satisfying the stated requirements. The corresponding ways in which one intensity is transferred into the other one are exactly the *most* and the *least* energy efficient in the sense of the Monge-Kantorovich cost. This result is thus ultimately a variant of Fermat's principle.

The fact that the solution to the above geometrical optics problem can be derived from a variational principle gives rise to a numerical treatment of the problem different from the one used in [11]; see also [12], [13]. In [11] the Monge-Ampère equation corresponding to (1.1) (see equation (2.6) in Section 2) was solved directly by a special geometric approximation by equations in measures with point-concentrated densities approximating its right-hand side. This method requires an iterative solution of a system of quadratic equations and involves frequent constructions of convex hulls in space. While the method is proved to converge [12], certain difficulties arise when the number of nodes becomes large. In the variational approach, when the problem of minimizing (1.2) under constraints (1.3) is discretized, we have a linear programming problem. However, in order to get a good approximation, one also has to deal here with a very large number of constraints and the issues of convergence and accuracy are open. We ran some numerical experiments with this approach and intend to return to this point in a separate publication. In this connection we also point out the work by Benamou and Brenier [2] in which the problem of finding the optimal solution to the Monge-Kantorovich problem with quadratic cost is solved numerically by transforming it into a special time dependent flow.

This paper is organized as follows. In Section 2, we recall some results from [13] concerning the ray tracing map, assuming smoothness of the reflectors, and formulate the main “two-reflector” problem. In Section 3, we give a geometric characterization of reflectors as envelopes of certain families of paraboloids. Such a characterization is of independent interest. In Section 4 we use this geometric characterization to define weak solutions of type A and type B of the two-reflector problem. To prove existence and uniqueness of solutions for each type we utilize the ideas of the Monge-Kantorovich theory and introduce the functional (1.2) on a certain class of “quasi-reflector” systems. This is done in Section 5. In the same section it is shown that the problem of finding weak solutions of type A is equivalent to finding minimizers of (1.2). Weak solutions of type B correspond to maximizers of (1.2). On the other hand, existence of minimizers (maximizers) to this functional is not difficult and has been established before in [3], [6], [7]. This implies existence of solutions. Uniqueness in each of the respective classes is established in Section 6 by proving that the ray tracing map \tilde{P} associated with a weak solution minimizes or maximizes the quadratic Monge-Kantorovich cost for which the functional (1.2) is the dual. The main theorem on existence and uniqueness of weak solutions to the two-reflector problem is stated and proved in Section 7.

Finally, we note that similar methods can be used to formulate and solve other geometrical optics problems involving systems with single and multiple reflectors¹.

¹ **Added in proof** (January 27, 2004). Variational treatments of the problem of designing an optical system with a single reflector transforming the energy of a point source into a prespecified energy distribution on the far-sphere were given by Glimm-Oliker in [8] and by Wang in [16].

2. STATEMENT OF THE PROBLEM

We begin by reviewing briefly the analytic formulation of the problem for smooth reflectors; see [13] for more details.

Let R_1 be given by the position vector $r_1(x) = (x, z(x))$, $x \in \bar{\Omega}$, with $z \in C^2(\bar{\Omega})$. The unit normal u on R_1 is given by

$$u = \frac{(-\nabla z, 1)}{\sqrt{1 + |\nabla z|^2}}.$$

Consider a ray labeled by $x \in \bar{\Omega}$ and propagating in the positive direction k of the z -axis. According to the reflection law the direction of the ray $y(x)$ reflected off R_1 is given by

$$y = k - 2\langle k, u \rangle u = k - 2 \frac{(-\nabla z, 1)}{1 + |\nabla z|^2},$$

where \langle, \rangle is the inner product in R^{n+1} . Denote by $t(x)$ the distance from reflector R_1 to reflector R_2 along the ray reflected in the direction $y(x)$ and let $s(x)$ be the distance from R_2 to the wavefront \tilde{T}_d along the corresponding ray reflected off R_2 . Assume for now that $t \in C^1(\bar{\Omega})$ and R_2 is a C^1 hypersurface. The total optical path length (OPL) corresponding to the ray associated with the point $x \in \bar{\Omega}$ is $\ell(x) = z(x) + t(x) + s(x)$. A calculation shows that $\ell(x) = \text{const} \equiv \ell$ on $\bar{\Omega}$. Since

$$(2.1) \quad R_2 : \quad r_2(x) = r_1(x) + t(x)y(x), \quad x \in \bar{\Omega},$$

the image of x on the reflected wavefront \tilde{T}_d is given by

$$(2.2) \quad P_d(x) = r_1(x) + t(x)y(x) + s(x)k, \quad x \in \bar{\Omega}.$$

The equation (2.2) is the ray tracing equation for this two-reflector system.

Introduce the map $P(x) = P_d(x) - dk : \bar{\Omega} \rightarrow \tilde{T}$. A calculation [13] shows that

$$(2.3) \quad p = P(x) = x + \beta \nabla z(x), \quad x \in \bar{\Omega},$$

where $\beta = \ell - d$ is the “reduced” optical path length.

To simplify the notation we will write $L(P(x))$ instead of $L(P_d(x))$ ($\equiv L(P(x) + dk$). For the input intensity $I(x)$, $x \in \Omega$, and the output intensity $L(P(x))$ on T_d we have, in accordance with the differential form of the energy conservation law (1.1),

$$(2.4) \quad L(P(x))|J(P(x))| = I(x), \quad x \in \Omega,$$

where we also take into account that $J(P_d) = J(P)$. It follows from (2.4) that Ω , T , I , and L must satisfy the necessary condition

$$(2.5) \quad \int_T L(p) \, dp = \int_\Omega I(x) \, dx.$$

It follows from (2.3) that $J(P) = \det[\text{Id} + \beta \text{Hess}(z)]$, where Id is the identity matrix and Hess is the Hessian. Hence, by (2.4),

$$(2.6) \quad L(x + \beta \nabla z) | \det[\text{Id} + \beta \text{Hess}(z)] | = I, \quad \text{in } \Omega.$$

Thus, the problem of determining the reflectors R_1 and R_2 with properties described in the introduction requires solving the following problem.

Problem I. *Given bounded domains Ω and T on the hyperplane α and two nonnegative, integrable functions I on Ω and L on T satisfying (2.5), it is required to find a function $z \in C^2(\bar{\Omega})$ such that the map*

$$(2.7) \quad P_\alpha = x + \beta \nabla z : \bar{\Omega} \rightarrow \bar{T}$$

is a diffeomorphism satisfying equation (2.6).

It is shown in [11] that once such a function z is found, the function w describing the second reflector is determined by z and β as

$$(2.8) \quad w(P(x)) = d - s(x) = z(x) + \frac{\beta}{2} (|\nabla z|^2 - 1).$$

Following [13] we introduce the function

$$(2.9) \quad V(x) = \frac{x^2}{2} + \beta z(x) - \frac{\beta^2}{2}.$$

Then by (2.3) and (2.8)

$$(2.10) \quad P = \nabla V,$$

$$(2.11) \quad w = \frac{1}{\beta} \left[V - \langle x, \nabla V \rangle + \frac{1}{2} |\nabla V|^2 \right],$$

where $V - \langle x, \nabla V \rangle$ is the negative of the usual Legendre transform of V . Thus, V is a potential for the map $P : \bar{\Omega} \rightarrow \bar{T}$. If P is a diffeomorphism, then the inverse of the transformation $(x, z(x)) \rightarrow (p, w(p))$, where $p = P(x)$, is given by

$$(2.12) \quad x(p) = P^{-1}(p) = p - \beta \nabla_p w(p),$$

$$(2.13) \quad z(p) = w(p) - \left(\frac{\beta}{2} |\nabla_p w(p)|^2 + \frac{\beta}{2} \right), \quad p \in \bar{T}.$$

In terms of the potential V the equation (2.6) becomes

$$(2.14) \quad L(\nabla V) | \det \text{Hess}(V) | = I \quad \text{in } \Omega,$$

which is an equation of Monge-Ampère type.

In order to clarify the relations between the parameters ℓ , d and β , note first that it is the reduced optical path length β that is intrinsic to the problem. The choice of cross sections of the fronts (that is, the selection of a particular value for d) is extraneous. In fact, it is easy to see that, if (z, w) are two-reflectors as above and we change d to d' , then $\ell' - d' = \ell - d = \beta$ and (z, w) are not affected by such change.

Finally, we note that the two-reflector system described above has the following two symmetries. For $\lambda \in \mathbb{R}^+$ put $z'(x) = \lambda z(x)$ and $\beta' = (1/\lambda)\beta$. Then

$$\begin{aligned} P'(x) &= P(x), \\ w'(p) &= \lambda w(p) + \frac{\beta}{2} \left(\lambda - \frac{1}{\lambda} \right), \end{aligned}$$

where $P'(x) = x + \beta' \nabla z'(x)$. In other words, the system is invariant under some combination of flattening (stretching) the first and translating and flattening (stretching) the second reflector.

Note also that a horizontal translation of both reflectors, that is, adding the same constant to z and w , does not change β and the map P .

3. GEOMETRIC CHARACTERIZATION OF REFLECTORS

We examine first more closely the relationship between the functions z , V and w for smooth reflectors. *For the rest of the paper we assume that Ω and T are bounded domains on the hyperplane α .* We continue to assume that the map $P : \bar{\Omega} \rightarrow \bar{T}$ is a diffeomorphism. Let $x \in \bar{\Omega}$, $p \in \bar{T}$, and

$$Q(x, p) = \langle x, p \rangle + \beta w(p) - \frac{p^2}{2}.$$

If $p = P(x)$, then by (2.9)–(2.11) we have $V(x) = Q(x, P(x))$. Denote by S_V the graph of V over $\bar{\Omega}$. Let $x_0 \in \bar{\Omega}$ and $p_0 = P(x_0)$. The tangent hyperplane to S_V at $(x_0, V(x_0))$ is given by the equation

$$Z = \langle x, p_0 \rangle - \langle x_0, p_0 \rangle + V(x_0),$$

where (x, Z) denotes an arbitrary point on that hyperplane. Taking into account (2.11), we obtain

$$(3.1) \quad \langle x, p_0 \rangle - \langle x_0, p_0 \rangle + V(x_0) = \langle x, p_0 \rangle + \beta w(p_0) - \frac{p_0^2}{2} = Q(x, p_0).$$

Since $V(x_0) = Q(x_0, p_0)$, we conclude that $Z = Q(x, p_0)$ is the tangent hyperplane to S_V at $(x_0, V(x_0))$. Consequently, if V is convex, then $Z = Q(x, p_0)$ is a supporting hyperplane to S_V from below, and if V is concave, then $Z = Q(x, p_0)$ is supporting to S_V from above (relative to positive direction of the z -axis). Because \bar{T} is bounded, there are no vertical tangent hyperplanes to the graph of V and we have

$$\begin{aligned} V(x) &\geq Q(x, p_0) \quad \text{for all } x \in \bar{\Omega}, \text{ if } V \text{ is convex,} \\ V(x) &\leq Q(x, p_0) \quad \text{for all } x \in \bar{\Omega}, \text{ if } V \text{ is concave.} \end{aligned}$$

Since for every $p \in \bar{T}$ the hyperplane $Q(x', p)$ is supporting to S_V at some $(x', V(x'))$, we get

$$(3.2) \quad V(x) \geq Q(x, p) \quad \text{for all } x \in \bar{\Omega}, p \in \bar{T}, \text{ if } V \text{ is convex,}$$

$$(3.3) \quad V(x) \leq Q(x, p) \quad \text{for all } x \in \bar{\Omega}, p \in \bar{T}, \text{ if } V \text{ is concave,}$$

and in both cases we have equalities if $p = P(x)$.

Let

$$U(p) = \frac{p^2}{2} - \beta w(p) - \frac{\beta^2}{2}, \quad R(x, p) = \langle x, p \rangle - \beta z(x) - \frac{x^2}{2}.$$

It follows from (3.2) and (3.3) that

$$(3.4) \quad U(p) \geq R(x, p) \quad \text{for all } x \in \bar{\Omega}, p \in \bar{T}, \text{ if } V \text{ is convex,}$$

$$(3.5) \quad U(p) \leq R(x, p) \quad \text{for all } x \in \bar{\Omega}, p \in \bar{T}, \text{ if } V \text{ is concave,}$$

and in both cases equalities are achieved if $p = P(x)$. Also, for any fixed $x_0 \in \bar{\Omega}$ the hyperplane $R(x_0, p)$ is supporting to the graph S_U of U at $(p_0 = P(x_0), U(p_0))$.

Using the usual characterization of convex functions [15] we obtain from (3.2), (3.4) and (3.3), (3.5)

$$\begin{aligned} V(x) &= \sup_{p \in \bar{T}} Q(x, p), \quad U(p) = \sup_{x \in \bar{\Omega}} R(x, p) \quad \text{when } V \text{ is convex,} \\ V(x) &= \inf_{p \in \bar{T}} Q(x, p), \quad U(p) = \inf_{x \in \bar{\Omega}} R(x, p) \quad \text{when } V \text{ is concave.} \end{aligned}$$

For convex V this implies

$$(3.6) \quad z(x) = \sup_{p \in \bar{T}} \left[\frac{\beta^2 - |x - p|^2}{2\beta} + w(p) \right], \quad x \in \bar{\Omega},$$

$$(3.7) \quad w(p) = \inf_{x \in \bar{\Omega}} \left[\frac{|x - p|^2 - \beta^2}{2\beta} + z(x) \right], \quad p \in \bar{T}.$$

Similarly, when V is concave we have

$$(3.8) \quad z(x) = \inf_{p \in \tilde{T}} \left[\frac{\beta^2 - |x - p|^2}{2\beta} + w(p) \right], \quad x \in \tilde{\Omega},$$

$$(3.9) \quad w(p) = \sup_{x \in \tilde{\Omega}} \left[\frac{|x - p|^2 - \beta^2}{2\beta} + z(x) \right], \quad p \in \tilde{T}.$$

The characterizations (3.6), (3.7) and (3.8), (3.9) have a simple geometric meaning. To describe it, consider first the case when V is convex. Recall that the total optical path length $\ell = z(x) + t(x) + d - w(P(x)) = \text{const}$ (see Figure 1). Also, $t^2(x) = |x - P(x)|^2 + |z(x) - w(P(x))|^2$ and $\beta = \ell - d$. It follows from (3.6) that $(x, z(x))$ is a point on the graph of the paraboloid

$$(3.10) \quad k_{p,w}(x) = \frac{\beta^2 - |x - p|^2}{2\beta} + w, \quad x \in \alpha,$$

with the focus at $(p = P(x), w(P(x)))$ and focal parameter β .

Similarly, it follows from (3.7) that a point $(p = P(x), w(P(x)))$ on the second reflector lies on a paraboloid

$$(3.11) \quad h_{x,z}(p) = \frac{|x - p|^2 - \beta^2}{2\beta} + z, \quad p \in \alpha,$$

with the focus at $(x, z(x))$ and focal parameter β .

Let $K_{p,w(p)}$ be the convex body bounded by the graph of paraboloid $k_{p,w}(x)$, and $H_{x,z(x)}$ the convex body bounded by the graph of paraboloid $h_{x,z}(p)$. Then (3.6) and (3.7) mean that the graphs S_z of $z(x)$ and S_w of $w(p)$ are given by

$$(3.12) \quad S_z = \partial \left(\bigcup_{p \in \tilde{T}} K_{p,w(p)} \right),$$

$$(3.13) \quad S_w = \partial \left(\bigcup_{x \in \tilde{\Omega}} H_{x,z(x)} \right).$$

When the potential V is concave we have similar characterizations of S_z and S_w , with \bigcup in (3.12), (3.13) replaced by \bigcap .

Remark 3.1. It follows from (3.6), (3.7) that when $V(x)$ is convex, then for any $x \in \Omega$ the path taken by the light ray through the reflector system is the shortest among all possible paths (not necessarily satisfying the reflection law) that go from $(x, 0)$ to $(x, z(x))$, then to some $(p, w(p))$ and then to (p, d) . Of course, the shortest path satisfies the reflection law and $p = P(x)$. For concave V the corresponding light path is the longest as it follows from (3.8), (3.9). Thus the characterizations (3.6), (3.7) and (3.8), (3.9) are variants of the Fermat principle.

4. WEAK SOLUTIONS OF PROBLEM I

We use the geometric characterizations of reflectors in Section 3 to define weak solutions to Problem I. Let Ω and T be two bounded domains on the hyperplane α , and β a fixed positive number.

Definition 4.1. A pair $(z, w) \in C(\bar{\Omega}) \times C(\bar{T})$ is called a *two-reflector of type A* if

$$(4.1) \quad z(x) = \sup_{p \in \bar{T}} k_{p,w(p)}(x), \quad x \in \bar{\Omega},$$

$$(4.2) \quad w(p) = \inf_{x \in \bar{\Omega}} h_{x,z(x)}(p), \quad p \in \bar{T},$$

where $k_{p,w(p)}(x)$ and $h_{x,z(x)}(p)$ are defined by (3.10) and (3.11). Similarly, a pair $(z, w) \in C(\bar{\Omega}) \times C(\bar{T})$ is called a *two-reflector of type B* if

$$(4.3) \quad z(x) = \inf_{p \in \bar{T}} k_{p,w(p)}(x), \quad x \in \bar{\Omega},$$

$$(4.4) \quad w(p) = \sup_{x \in \bar{\Omega}} h_{x,z(x)}(p), \quad p \in \bar{T}.$$

To avoid repetitions, we consider below only two-reflectors of type A. The changes that need to be made to deal with two-reflectors of type B are straightforward and are omitted.

It will be convenient to construct the following extensions z^* of the function z , and w^* of w to the entire α . For a pair (z, w) as in definition 4.1 let

$$(4.5) \quad V(x) = \frac{x^2}{2} + \beta z(x) - \frac{\beta^2}{2} \quad \text{and} \quad Q(x, p) = \langle x, p \rangle + \beta w(p) - \frac{p^2}{2}.$$

It follows from (4.1) that

$$V(x) = \sup_{p \in \bar{T}} Q(x, p), \quad x \in \bar{\Omega}.$$

That is, V is convex and continuous over $\bar{\Omega}$. Furthermore, since \bar{T} is bounded, the graph S_V has no vertical supporting hyperplanes. For any fixed $p \in \bar{T}$ define the half-space $Q^+(p) = \{(x, Z) \in \alpha \times \mathbb{R}^1 \mid Z \geq Q(x, p)\}$. Then

$$S_{V^*} = \partial \left(\bigcap_{p \in \bar{T}} Q^+(p) \right)$$

is a graph of a convex function V^* defined for all $x \in \alpha$. Note that $V^*(x) = V(x)$ when $x \in \bar{\Omega}$. We now define an extension of z by putting

$$z^*(x) = \frac{1}{\beta} \left[-\frac{x^2}{2} + V^*(x) + \frac{\beta^2}{2} \right].$$

Similarly, for any fixed $x \in \bar{\Omega}$ we let $R^+(x) = \{(p, Z) \in \alpha \times \mathbb{R}^1 \mid Z \geq R(x, p)\}$, where

$$R(x, p) = \langle x, p \rangle - \beta z(x) - \frac{x^2}{2}, \quad p \in \alpha.$$

Then the function

$$U^*(p) = \sup_{x \in \bar{\Omega}} R(x, p), \quad p \in \alpha,$$

is defined. It is also convex. It follows from (4.2) that for $p \in \bar{T}$

$$U^*(p) = \frac{p^2}{2} - \beta w(p) - \frac{\beta^2}{2} \quad (\equiv U(p)).$$

The corresponding extension of w we define as

$$w^*(p) = \frac{1}{\beta} \left[\frac{p^2}{2} - U^*(p) - \frac{\beta^2}{2} \right], \quad p \in \alpha.$$

Lemma 4.2. *The function V^* is uniformly Lipschitz on α with Lipschitz constant $\max_{\bar{T}} |p|$. Also, $U^*(p)$, $p \in \alpha$, is uniformly Lipschitz on α with Lipschitz constant $\max_{\bar{\Omega}} |x|$. In addition, $z \in \text{Lip}(\bar{\Omega})$ and $w \in \text{Lip}(\bar{T})$, with the Lipschitz constant $\leq \sup_{(x,p) \in \bar{\Omega} \times \bar{T}} |x - p|/\beta$.*

Proof. By our convention the normal vector to a plane $Q(x, p)$ (when p is fixed) is given by $(-p, 1)$. It follows from definition of V^* that S_{V^*} has no supporting hyperplanes with normal $(-p, 1)$ such that $p \notin \bar{T}$. Since T is bounded, this implies the first statement of the lemma. The statements regarding U^* are established by similar arguments. From these properties of V^* it follows that the function z^* is continuous on α and Lipschitz on any compact subset of α . Similar properties hold also for w^* .

Now we estimate the Lipschitz constant for z on $\bar{\Omega}$. Let (z, w) be a two-reflector of type A. Let $x, x' \in \bar{\Omega}$ and let $z(x') \geq z(x)$. (If the opposite inequality holds we relabel x and x' .) Fix some small $\varepsilon > 0$. It follows from (4.1), (4.2) that there exists a $p' \in \bar{T}$ such that $z(x') \leq k_{p', w(p')}(x') + \varepsilon$. Then

$$\begin{aligned} |z(x') - z(x)| &\leq k_{p', w(p')}(x') - z(x) + \varepsilon \leq k_{p', w(p')}(x') - k_{p', w(p')}(x) + \varepsilon \\ &\leq \sup_{x \in \bar{\Omega}} |\nabla k_{p', w(p')}(x)| |x' - x| + \varepsilon = \frac{1}{\beta} \sup_{s \in \bar{\Omega}} |s - p'| |x' - x| + \varepsilon \\ &\leq \frac{1}{\beta} \sup_{s \in \bar{\Omega}, p \in \bar{T}} |s - p| |x' - x| + \varepsilon. \end{aligned}$$

Letting $\varepsilon \rightarrow 0$, we obtain the statement regarding the Lipschitz constant for z . The same statement regarding w and two-reflectors of type B are proved similarly. □

Next, we define *the analogue of the ray tracing map P* for a two-reflector. For that we need to recall the notion of the *normal map* [1], p. 114. Let $u : G \rightarrow \mathbb{R}^1$ be an arbitrary convex function defined on a domain $G \subset \alpha$, and S_u its graph. For $x_0 \in G$ let $Z - u(x_0) = \langle p, x - x_0 \rangle$ be a hyperplane with normal $(-p, 1)$ supporting to S_u at $(x_0, u(x_0))$. The normal map $v_u : G \rightarrow \alpha$ at x_0 is defined as

$$v_u(x_0) = \bigcup \{p\},$$

where the union is taken over all hyperplanes supporting to S_u at $(x_0, u(x_0))$.

Definition 4.3. Let (z, w) be a two-reflector of type A. For $x \in \alpha$ we put

$$\tilde{P}(x) = v_{V^*}(x).$$

For reflectors of type B the ray tracing map is defined similarly, using the function $-U^*$. In general, \tilde{P} may be multivalued.

Lemma 4.4. *Let (z, w) be a two-reflector of type A and z^* and w^* their respective extensions, as above. Then $\tilde{P}(x) \in \tilde{T}$ for all $x \in \alpha$. In addition, for any $p \in \tilde{T}$ the set $\{x \in \tilde{\Omega} \mid \tilde{P}(x) = p\} \neq \emptyset$. Furthermore, for any $x \in \tilde{\Omega}$*

$$(4.6) \quad \tilde{P}(x) = \{\text{all } p \in \tilde{T} \mid w(p) = h_{x,z(x)}(p)\}.$$

Proof. Let $x \in \alpha$ and $Q(x, p)$ a supporting hyperplane to V^* at $(x, V^*(x))$. Then the normal p is in $v_{V^*}(x)$. On the other hand, by definition of V^* , S_{V^*} has only supporting hyperplanes with normals in \tilde{T} . Hence, $\tilde{P}(x) \subset \tilde{T}$.

Let $p \in \tilde{T}$, and $Q(x, p)$ a supporting hyperplane to S_{V^*} . We need to show that there is an $x \in \tilde{\Omega}$ such that $p \in \tilde{P}(x)$. By (4.5) and (4.2) we have for any $x \in \tilde{\Omega}$

$$V(x) - Q(x, p) = \beta(h_{x,z(x)}(p) - w(p)) \geq 0.$$

By (4.2) there exists an $x \in \tilde{\Omega}$ such that $V(x) - Q(x, p) = 0$. This implies the remaining two statements of the lemma. □

Remark 4.5. It follows from definition that \tilde{P} is multivalued at points x where S_{V^*} has more than one supporting hyperplane. At such x the function z^* is not differentiable. Let $(x_0, z^*(x_0))$ be one such point. Then

$$\tilde{P}(x_0) = \{p \in \alpha \mid Q(x, p) \text{ is supporting to } S_{V^*} \text{ at } (x_0, V^*(x_0))\}.$$

In other words, a light ray labeled by $x_0 \in \Omega$ that hits a point where the first reflector has a singular point will split into a cone of light rays. These rays will generate a subset on the paraboloid $h_{x_0,z(x_0)}(p)$ whose projection on α is $\tilde{P}(x_0)$. This is consistent with the physical interpretation of diffraction at singularities of this type [9].

Remark 4.6. Since V^* is convex, by Rademacher's theorem, the Lebesgue measure of the set of singular points on S_{V^*} is zero. Thus, $\tilde{P}(x)$ is single-valued almost everywhere in α . Furthermore, the function z (z^*) is a difference of two convex functions, and therefore, it is differentiable almost everywhere in Ω (α). The same is true for w and w^* . It follows then from the definitions of \tilde{P} and V that for almost all $x \in \Omega$

$$(4.7) \quad \tilde{P}(x) = \nabla V = x + \beta \nabla z(x).$$

A similar property holds also for the function w .

Lemma 4.7. *If (z, w) is a two-reflector of type A, then for all $x \in \tilde{\Omega}$, $p \in \tilde{T}$*

$$(4.8) \quad z(x) - w(p) \geq \frac{1}{2\beta}(\beta^2 - |x - p|^2).$$

In addition, for almost all $x \in \Omega$ there exists a unique $p \in \tilde{T}$ such that $p = \tilde{P}(x)$, and (4.8) in this case is an equality.

Proof. The lemma follows from (4.1), (4.2), Remark 4.6, and Lemma 4.4. \square

Define the inverse of \tilde{P} for $p \in \tilde{T}$ as

$$\tilde{P}^{-1}(p) = \{x \in \alpha \mid p \in \tilde{P}(x)\}.$$

Theorem 4.8. *Let \mathcal{B} be the σ -algebra of Borel subsets of T . Let (z, w) be a two-reflector of type A. For any set $\tau \in \mathcal{B}$ the set $\tilde{P}^{-1}(\tau)$ is measurable relative to the standard Lebesgue measure on α . In addition, for any non-negative locally integrable function I on α the function*

$$\mathcal{L}(\tau) = \int_{\tilde{P}^{-1}(\tau)} I(x) dx$$

is a non-negative completely additive measure on \mathcal{B} .

Proof. The proof of this theorem is completely analogous to the proofs of Theorems 9 and 16 in [14]. \square

Lemma 4.9. *Let Ω and T be two bounded domains on α and I a non-negative integrable function on Ω extended to entire α by setting $I(x) \equiv 0$ for $x \in \alpha \setminus \Omega$. Let (z, w) be a two-reflector of type A or B. Then for any continuous function h on \tilde{T} we have the following change of variable formula*

$$(4.9) \quad \int_T h(p) \mathcal{L}(dp) = \int_{\Omega} h(\tilde{P}(x)) I(x) dx.$$

Proof. In the integral on the right $h(\tilde{P}(x))$ is discontinuous only where \tilde{P} is not single valued, that is, on the set of measure zero. Thus, the integral on the right is well defined.

We may assume that $\int_{\Omega} I(x) dx > 0$; otherwise, the statement is trivial. Fix some small $\varepsilon > 0$ and a positive integer N . Partition the interval $[\min h(p), \max h(p)]$ into sub-intervals S_1, \dots, S_N of length $< \varepsilon / \int_{\Omega} I(x) dx$ and let $h_i \in S_i$. Put $\tau_i = \{p \in T \mid h(p) \in S_i\}$. Then for sufficiently large N

$$\left| \int_T h(p) \mathcal{L}(dp) - \sum h_i \mathcal{L}(\tau_i) \right| < \varepsilon.$$

For any $i, j = 1, \dots, N, i \neq j$, $\text{meas}(\tilde{P}^{-1}(\tau_i) \cap \tilde{P}^{-1}(\tau_j)) = 0$ (see Remark 4.6). Hence,

$$\left| \int_{\Omega} h(\tilde{P}(x)) I(x) dx - \sum h_i \int_{\tilde{P}^{-1}(\tau_i)} I(x) dx \right| < \varepsilon.$$

This, together with the previous inequality, imply

$$\left| \int_{\Omega} h(\tilde{P}(x)) I(x) dx - \sum h_i \mathcal{L}(\tau_i) \right| < 2\varepsilon. \quad \square$$

Definition 4.10. A two-reflector (z, w) of type A (B) is called a weak solution of type A (B) of the two-reflector Problem I if the map $\tilde{P} : \tilde{\Omega} \rightarrow \tilde{T}$ is onto, and for any Borel set $\tau \subseteq T$

$$\mathcal{L}(\tau) = \int_{\tau} L(p) dp.$$

Using Lemma 4.9 and this definition we obtain the following result.

Lemma 4.11. *Let (z, w) be a weak solution of type A (B) of the two-reflector problem I. Then for any continuous function h on \tilde{T}*

$$(4.10) \quad \int_T h(p) L(p) dp = \int_{\Omega} h(\tilde{P}(x)) I(x) dx.$$

5. A VARIATIONAL PROBLEM AND WEAK SOLUTIONS OF THE TWO-REFLECTOR PROBLEM

As before, we consider here only the case of two-reflectors of type A. We comment on the case of two-reflectors of type B in Section 7. Let, as before, ℓ and d be the given parameters of the system, and $\beta = \ell - d$. Let $\zeta \in C(\tilde{\Omega})$ and $\omega \in C(\tilde{T})$. With any such pair (ζ, ω) we associate a “quasi-reflector” system in which the light path is defined as follows. Let $(x, p) \in \tilde{\Omega} \times \tilde{T}$ and let $(x, \zeta(x))$ be the point where the horizontal ray emanating from $(x, 0)$ hits the graph of ζ . Let P_2 be the

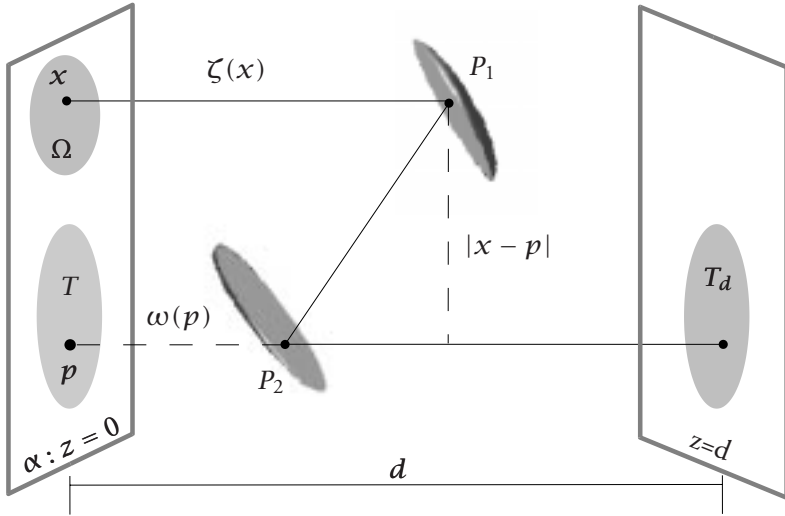


FIGURE 2. Definition of $\ell(\zeta, \omega, x, p)$.

point where the horizontal ray begins in order to terminate in (p, d) . The “light” path is the polygon $(x, 0)P_1P_2(p, d)$; see Figure 2. The length of this path is

$$\ell(\zeta, \omega, x, p) = \zeta(x) + \sqrt{(\zeta(x) - \omega(p))^2 + |x - p|^2} + d - \omega(p).$$

The class of *admissible pairs* is defined as

$$(5.1) \quad \text{Adm}(\Omega, T) = \{(\zeta, \omega) \in C(\bar{\Omega}) \times C(\bar{T}) \mid \ell(\zeta, \omega, x, p) \geq \ell, \forall (x, p) \in \bar{\Omega} \times \bar{T}\}.$$

By construction, a pair $(\zeta, \omega) \in C(\bar{\Omega}) \times C(\bar{T})$ lies in $\text{Adm}(\Omega, T)$ if and only if for all $x \in \bar{\Omega}, p \in \bar{T}$

$$(5.2) \quad \zeta(x) \geq k_{p, \omega(p)}(x), \quad \omega(p) \leq h_{x, \zeta(x)}(p),$$

or, equivalently, if and only if

$$(5.3) \quad \zeta(x) - \omega(p) \geq \frac{1}{2\beta}(\beta^2 - |x - p|^2),$$

for all $x \in \bar{\Omega}, p \in \bar{T}$. It follows from (4.1), (4.2) that a two-reflector of type A is an admissible pair.

The following functional is central to our investigation. Let I and L be two non-negative integrable functions on Ω and T , respectively, satisfying the energy conservation law (2.5). For $(\zeta, \omega) \in \text{Adm}(\Omega, T)$ put

$$\mathcal{F}(\zeta, \omega) = \int_{\Omega} \zeta(x)I \, dx - \int_T \omega(p)L \, dp.$$

Clearly, \mathcal{F} is linear and bounded on $C(\bar{\Omega}) \times C(\bar{T})$ with respect to the norm $\max\{\|\zeta\|_{\infty}, \|\omega\|_{\infty}\}$. Geometrically, $\mathcal{F}(\zeta, \omega)$ is proportional to the mean horizontal distance between the points of the two graphs, the average being weighted by the intensities.

We consider now the following problem.

Problem II. Minimize \mathcal{F} on $\text{Adm}(\Omega, T)$.

Proposition 5.1 ([3, 6, 7]). There exist $(z, w) \in \text{Adm}(\Omega, T)$ such that

$$\mathcal{F}(z, w) = \inf_{(\zeta, \omega) \in \text{Adm}(\Omega, T)} \mathcal{F}(\zeta, \omega).$$

We may further assume that the pair (z, w) satisfies the conditions (4.1), (4.2) for a type A reflector system.

Proof. It follows from (5.2) that one can restrict the search for minimizers to two-reflectors (ζ, ω) of type A. Also, because of (2.5) $\mathcal{F}(\zeta, \omega)$ is invariant under translations $\zeta \mapsto \zeta + \rho$, $\omega \mapsto \omega + \rho$ for a constant $\rho \in \mathbb{R}$. Hence, it is sufficient to consider only two-reflectors (ζ, ω) for which $\zeta(x_0) = 0$ for some $x_0 \in \bar{\Omega}$.

By Lemma 4.2, ζ and ω are uniformly Lipschitz, with the Lipschitz constant $K = \sup_{x \in \Omega, p \in T} |x - p|/\beta$. It follows that for all $x \in \Omega$,

$$|\zeta(x)| = |\zeta(x) - \zeta(x_0)| \leq K \text{diam } \Omega.$$

For all $p \in \bar{T}$ we have

$$\omega(p) \leq h_{x_0, \zeta(x_0)}(p) = \frac{1}{2\beta}(|x_0 - p|^2 - \beta^2) \leq \max_{q \in \bar{T}} \frac{1}{2\beta}(|x_0 - q|^2 - \beta^2).$$

Finally, since $h_{x, \zeta(x)}(p) = (|x - p|^2 - \beta^2)/2\beta + \zeta(x) \geq -\beta/2 - K \text{diam } \Omega$ for all $x \in \bar{\Omega}$, we also get the lower bound

$$\omega(p) = \inf_{x \in \Omega} h_{x, \zeta(x)}(p) \geq -\frac{\beta}{2} - K \text{diam } \Omega.$$

Therefore, the maps ζ and ω are also uniformly bounded. By the Arzelà-Ascoli theorem, and because \mathcal{F} is continuous, the infimum of \mathcal{F} is achieved at some two-reflector pair (z, w) . □

We now show that weak solutions of Problem I and solutions of Problem II are the same. This establishes the existence of weak solutions to Problem I.

Theorem 5.2. *Let $z \in C(\bar{\Omega})$, $w \in C(\bar{T})$ be a two-reflector of type A. Then the following statements are equivalent.*

- (i) (z, w) minimizes \mathcal{F} in $\text{Adm}(\Omega, T)$.
- (ii) (z, w) is a weak solution of type A of the two-reflector Problem I.

Proof. The proof below is similar to the proofs of Theorem 1 in [7] and of the change of variable formula in [6]. In order to make our presentation reasonably self-contained we present it here.

(ii) \Rightarrow (i). Let $(\zeta, \omega) \in \text{Adm}(\Omega, T)$. Then by (5.3) and Lemma 4.7, for almost all $x \in \Omega$,

$$\zeta(x) - \omega(\bar{P}(x)) \geq \frac{1}{2\beta}(\beta^2 - |x - \bar{P}(x)|^2) = z(x) - w(\bar{P}(x)).$$

Integrating this inequality we get

$$\int_{\Omega} \zeta(x)I dx - \int_{\Omega} \omega(\bar{P}(x))I dx \geq \int_{\Omega} z(x)I dx - \int_{\Omega} w(\bar{P}(x))I dx.$$

Using Lemma 4.11, we get

$$\begin{aligned} \mathcal{F}(\zeta, \omega) &= \int_{\Omega} \zeta(x)I dx - \int_T \omega(p)L dp \\ &\geq \int_{\Omega} z(x)I dx - \int_T w(p)L dx = \mathcal{F}(z, w). \end{aligned}$$

Since (ζ, ω) was arbitrary, we are done.

(i) \Rightarrow (ii). Let \bar{P} denote the ray tracing map for the pair (z, w) . It must be shown that $\int_{\bar{P}^{-1}(\tau)} I(x) dx = \int_{\tau} L(p) dp$ for all Borel sets $\tau \subseteq T$. We will prove that this is the Euler-Lagrange equation for the functional \mathcal{F} .

It is sufficient to establish this for the case when τ is an open ball with center $p_0 \in T$ and radius $r > 0$, contained in T . For $i = 1, 2, \dots$, define for $p \in \alpha$

$$\chi_i(p) = \begin{cases} 1, & \text{if } |p - p_0| < r - 1/i, \\ i(r - |p - p_0|), & \text{if } r - 1/i \leq |p - p_0| < r, \\ 0, & \text{if } |p - p_0| \geq r. \end{cases}$$

Then χ_i is continuous on α , $0 \leq \chi_i \leq 1$, and the sequence $\{\chi_i(p) \mid p \in \alpha\}_{i=1}^{\infty}$ converges on α pointwise to the characteristic function of τ , $\chi_{\tau}(p)$.

Fix some i and for $\varepsilon \in (-1, 1)$ put

$$w_\varepsilon(p) = w(p) + \varepsilon \cdot \chi_i(p).$$

$$z_\varepsilon(x) = \sup_{p \in \bar{T}} k_{p, w_\varepsilon(p)}(x) = \sup_{p \in \bar{T}} \left\{ \frac{1}{2\beta} (\beta^2 - |x - p|^2) + w_\varepsilon(p) \right\}.$$

By construction, the pair $(z_\varepsilon, w_\varepsilon)$ satisfies the condition (5.3), with ζ replaced by z_ε and ω replaced by w_ε . We show now that z_ε belongs to $\text{Lip}(\bar{\Omega})$. Let $x, x' \in \bar{\Omega}$ and let $z(x') \geq z(x)$. (If the opposite inequality holds, we relabel x and x' .) Let $p' \in \bar{T}$ be such that $z_\varepsilon(x') = k_{p', w_\varepsilon(p')}(x')$. Then

$$\begin{aligned} |z_\varepsilon(x') - z_\varepsilon(x)| &= k_{p', w_\varepsilon(p')}(x') - z_\varepsilon(x) \leq k_{p', w_\varepsilon(p')}(x') - k_{p', w_\varepsilon(p')}(x) \\ &\leq \sup_{s \in \bar{\Omega}} |\nabla k_{p', w_\varepsilon(p')}(s)| |x' - x| = \frac{1}{\beta} \sup_{s \in \bar{\Omega}} |s - p'| |x' - x| \\ &\leq \frac{1}{\beta} \sup_{s \in \bar{\Omega}, p \in \bar{T}} |s - p| |x' - x|. \end{aligned}$$

Hence, z_ε is continuous and $(z_\varepsilon, w_\varepsilon) \in \text{Adm}(\Omega, T)$.

Now let $x \in \bar{\Omega}$. For each ε let p_ε be a point in \bar{T} such that $z_\varepsilon(x) = k_{p_\varepsilon, w_\varepsilon(p_\varepsilon)}(x)$. (This choice, of course, may not be unique.) Then

$$\begin{aligned} z_\varepsilon(x) - z(x) &= k_{p_\varepsilon, w_\varepsilon(p_\varepsilon)}(x) - z(x) \\ &\leq k_{p_\varepsilon, w_\varepsilon(p_\varepsilon)}(x) - k_{p_\varepsilon, w}(p_\varepsilon)(x) \\ &= w_\varepsilon(p_\varepsilon) - w(p_\varepsilon) = \varepsilon \chi_i(p_\varepsilon). \end{aligned}$$

Similarly, if $p \in \bar{P}(x)$, then

$$z_\varepsilon(x) - z(x) = z_\varepsilon(x) - k_{p, w(p)}(x) \geq k_{p, w_\varepsilon(p)}(x) - k_{p, w(p)}(x) = \varepsilon \chi_i(p).$$

Therefore,

$$(5.4) \quad -|\varepsilon| \leq \varepsilon \chi_i(p) \leq z_\varepsilon(x) - z(x) \leq \varepsilon \chi_i(p_\varepsilon) \leq |\varepsilon|,$$

for all $x \in \bar{\Omega}$. In particular, z_ε converges uniformly to z on $\bar{\Omega}$, as $\varepsilon \rightarrow 0$.

Now consider those $x \in \Omega$ for which the ray tracing map \bar{P} is single-valued. This is the case for almost all $x \in \Omega$. We claim that $p_\varepsilon = \bar{P}_\varepsilon(x)$, where \bar{P}_ε denotes the ray tracing map for $(z_\varepsilon, w_\varepsilon)$, converges to p as $\varepsilon \rightarrow 0$.

Suppose this is not true. Then there is a sequence $\{p_{\varepsilon_j}\}$, $j = 1, 2, \dots$, with $\varepsilon_j \rightarrow 0$ as $j \rightarrow \infty$ and a constant $\eta > 0$ such that

$$|p - p_{\varepsilon_j}| > \eta, \quad \text{for all } j.$$

Let $z'(x) = \max_{p' \in \bar{T}, |p' - p| \geq \eta} k_{p', w}(p')(x)$. Note that the maximum in the definition of z' is attained. Since p is the unique point in \bar{T} such that $z(x) = k_{p, w}(p)(x)$, it follows that $z' < z(x)$. Therefore, for all j ,

$$\begin{aligned} z(x) - z_{\varepsilon_j}(x) &= z(x) - k_{p_{\varepsilon_j}, w_{\varepsilon_j}(p_{\varepsilon_j})}(x) \\ &= z(x) - k_{p_{\varepsilon_j}, w}(p_{\varepsilon_j})(x) + w(p_{\varepsilon_j}) - w_{\varepsilon_j}(p_{\varepsilon_j}) \\ &= z(x) - k_{p_{\varepsilon_j}, w}(p_{\varepsilon_j})(x) - \varepsilon_j \chi_i(p_{\varepsilon_j}) \geq z(x) - z' - |\varepsilon_j|. \end{aligned}$$

This contradicts the fact that $z_{\varepsilon_j}(x)$ converges to $z(x)$ as $j \rightarrow \infty$. Hence, $p_\varepsilon \rightarrow p$ if $\varepsilon \rightarrow 0$.

It follows from (5.4) that

$$\left| \frac{z_\varepsilon(x) - z(x)}{\varepsilon} - \chi_i(p) \right| \leq |\chi_i(p_\varepsilon) - \chi_i(p)|.$$

Letting $\varepsilon \rightarrow 0$ and using the continuity of χ_i , we conclude that for almost all $x \in \Omega$

$$\frac{d}{d\varepsilon} \Big|_{\varepsilon=0} z_\varepsilon(x) = \chi_i(p) = \chi_i(\bar{P}(x)).$$

Then

$$\frac{d}{d\varepsilon} \Big|_{\varepsilon=0} \int_\Omega z_\varepsilon(x) I(x) dx = \int_\Omega \chi_i(\bar{P}(x)) I(x) dx.$$

Since \mathcal{F} has a minimum at (z, w) , we obtain

$$0 = \frac{d}{d\varepsilon} \Big|_{\varepsilon=0} \mathcal{F}(z_\varepsilon, w_\varepsilon) = \int_\Omega \chi_i(\bar{P}(x)) I(x) dx - \int_T \chi_i(p) L(p) dp.$$

Now let $i \rightarrow \infty$ in this equality. This is possible as $\chi_i(p) \rightarrow \chi_\tau(p)$ pointwise on α and therefore $\chi_i(\bar{P}(x)) \rightarrow \chi_\tau(\bar{P}(x))$ pointwise almost everywhere on Ω . Then, noting that $\chi_\tau(\bar{P}(x)) = \chi_{\bar{P}^{-1}(\tau)}(x)$ for almost all $x \in \Omega$, we obtain

$$\begin{aligned} \int_\tau L(p) dp &= \int_T \chi_\tau(p) L(p) dp = \int_\Omega \chi_\tau(\bar{P}(x)) I(x) dx \\ &= \int_\Omega \chi_{\bar{P}^{-1}(\tau)}(x) I(x) dx = \int_{\bar{P}^{-1}(\tau)} I(x) dx. \end{aligned}$$

This completes the proof of the theorem. □

6. CONNECTION BETWEEN THE TWO-REFLECTOR AND MONGE-KANTOROVICH PROBLEMS

In our notation, the Monge-Kantorovich mass transfer problem [3] can be formulated as follows. Consider the class of maps $P : \Omega \rightarrow T$ which are measure-preserving, that is, they satisfy the substitution rule

$$\int_\Omega h(P(x)) I dx = \int_T h(p) L dp,$$

for all continuous functions h on \bar{T} . Each such map is called a *plan*.

Problem III. *Minimize the quadratic transportation cost*

$$(6.1) \quad P \mapsto \frac{1}{2} \int_{\Omega} |x - P(x)|^2 I \, dx.$$

among all planes P .

Note that, for weak solutions to the two-reflector Problem I, the ray tracing map \tilde{P} is a plan by Lemma 4.11. In fact, we have the following result.

Theorem 6.1. *Let (z, w) be a weak solution of type A of the two-reflector Problem I. Let \tilde{P} be the corresponding ray tracing map. Then \tilde{P} minimizes the quadratic transportation cost (6.1) among all planes $P : \Omega \rightarrow T$, and any other minimizer is equal to \tilde{P} almost everywhere on $\text{supp}(I) \setminus \{x \in \bar{\Omega} \mid I(x) = 0\}$.*

Proof. Let $P : \Omega \rightarrow T$ be any plan. Then by (4.8),

$$(6.2) \quad z(x) - w(P(x)) \geq \frac{1}{2\beta} (\beta^2 - |x - P(x)|^2),$$

for all $x \in \Omega$ and equality holds if and only if $P(x) = \tilde{P}(x)$ for almost all $x \in \Omega$. Integrating against $I \, dx$ and applying Lemma 4.11, we get

$$\begin{aligned} & \frac{1}{2\beta} \int_{\Omega} [\beta^2 - |x - P(x)|^2] I \, dx \\ & \leq \int_{\Omega} [z(x) - w(P(x))] I \, dx = \int_{\Omega} z(x) I \, dx - \int_T w(p) L \, dp \\ & = \int_{\Omega} [z(x) - w(\tilde{P}(x))] I \, dx = \frac{1}{2\beta} \int_{\Omega} [\beta^2 - |x - \tilde{P}(x)|^2] I \, dx. \end{aligned}$$

This shows that \tilde{P} is a minimizer of the transportation cost.

To show uniqueness, note that if equality holds in the integral inequality, then equality must hold in (6.2) for almost all $x \in \text{supp}(I) \setminus \{I = 0\}$. Therefore, $P \equiv \tilde{P}$ a.e. on $\text{supp}(I) \setminus \{I = 0\}$. □

Remark 6.2. The functional \mathcal{F} is the dual of the quadratic cost functional (6.1) [3].

7. EXISTENCE AND UNIQUENESS OF WEAK SOLUTIONS TO THE TWO-REFLECTOR PROBLEM

Theorem 7.1. *There exist weak solutions of type A and of type B to the two-reflector Problem I. If (z, w) is any such solution pair then $z \in \text{Lip}(\bar{\Omega})$ and $w \in$*

$\text{Lip}(\bar{T})$. The corresponding ray-tracing map is single-valued almost everywhere and for almost all $x \in \Omega$ it is given by

$$\bar{P} = x + \beta \nabla z(x).$$

Furthermore, if (z, w) and (z', w') are two solutions of the same type with ray tracing maps \bar{P} and \bar{P}' , respectively, then

$$\bar{P}(x) \equiv \bar{P}'(x),$$

for almost all $x \in \text{supp}(I) \setminus \{x \in \Omega \mid I(x) = 0\}$.

Proof. By Proposition 5.1 and Theorem 5.2 we know that Problem I has a solution. The property that $z \in \text{Lip}(\bar{\Omega})$, $w \in \text{Lip}(\bar{T})$ follows from Lemma 4.2 and Remark 4.6. The same remark implies the statement regarding the ray-tracing map. The only property that remains to be checked is that for that solution, $\bar{P}: \bar{\Omega} \rightarrow \bar{T}$ is onto. But this follows from Lemma 4.4.

In order to prove uniqueness, let (z, w) and (z', w') be two solutions of type A, with ray tracing maps \bar{P} and \bar{P}' , respectively. Then by Proposition 6.1, both are minimizers of the quadratic Monge-Kantorovich cost functional, so that $\bar{P} = \bar{P}'$ a.e. on $\text{supp}(I) \setminus \{I = 0\}$.

Remark 7.2. Further regularity of weak solutions to (2.6) under additional assumptions on domains Ω and T and the density functions I and L follows from the regularity theory of Caffarelli [5], [4].

So far, existence and uniqueness of weak solutions to the two-reflector Problem I has been shown for weak solutions of type A. However, it is clear that a similar result holds for weak solutions of type B as well. Namely, we can define $\text{Adm}_+(\Omega, T)$ as the space of all pairs of reflectors such that $\ell(\xi, \omega, x, p)$ is less than ℓ , and then that \mathcal{F} admits a maximizing pair on $\text{Adm}_+(\Omega, T)$, which is a weak solution of the two-reflector Problem I. This shows, in particular, that for such a solution the ray tracing map maximizes the quadratic transportation cost among all planes. □

Corollary 7.3. Suppose that, in addition to the assumptions in Theorem 7.1, the function $I > 0$ in Ω . Then there is a constant $\rho \in \mathbb{R}$ such that

$$\begin{aligned} z'(x) &= z(x) + \rho, \\ w'(p) &= w(p) + \rho, \end{aligned}$$

for all $x \in \bar{\Omega}$ and all $p \in \bar{T}$. In other words, weak solutions of type A are unique on $\bar{\Omega}$ and \bar{T} up to a translation of the reflector system, and the same result holds for type B solutions.

Proof. We show this for weak solutions of type A. By the theorem, $\tilde{P}(x) \equiv \tilde{P}'(x)$ for all $x \in \Omega$. Note that by (4.7), $\nabla z(x) = \nabla z'(x)$ for almost all $x \in \Omega$. It follows that there is a constant ρ such that $z'(x) = z(x) + \rho$ for all $x \in \bar{\Omega}$. Now, by definition of $w'(p)$,

$$w'(p) = \inf_{x \in \bar{\Omega}} h_{x,z'(x)}(p) = \inf_{x \in \bar{\Omega}} h_{x,z(x)}(p) + \rho = w(p) + \rho,$$

for all $p \in \bar{T}$. □

Acknowledgement. The research of the second author was partially supported by a grant from Emory University Research Committee and by a National Science Foundation grant DMS-04-05622.

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KEY WORDS AND PHRASES: geometrical optics, reflectors, PDE's of Monge-Ampère type, Monge-Kantorovich mass transfer problem, measure preserving maps.

2000 MATHEMATICS SUBJECT CLASSIFICATION: 35J65, 78A05, 49K20.

Received: April 13th, 2003; revised: July 15th, 2003.